High-efficiency organic polymer light-emitting heterostructure devices on flexible plastic substrates
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Flexible plastic substrate has a distinct advantage over glass substrate in many applications. Unlike glass, the plastic is usually more robust and compact, has lighter weight, and is more cost effective. The use of the plastic substrates will also enable new product concepts such as curved displays, smart cards, and all plastic electronics. Because of these advantages, in the last decade, plastic substrates have been used in supertwisted nematic liquid-crystal displays (STN-LCDs),1 active-matrix liquid-crystal displays (AMLCDs),2 and organic light-emitting devices (OLEDs).3,4 While it is a very promising new technology, the plastic substrate, however, is limited by its low processing temperature, high heat-induced shrinkage, high gas permeability, low chemical resistance, average water and solvent durability, and average transparency. The low thermal durability of the plastic substrate prevents, for example, the high temperature postdeposition thermal annealing of the transparent conducting oxide (TCO), which is typically used in the organic light-emitting device as the transparent electrode. As a result, the TCO on the plastic substrate usually has a lower transparency, a lower conductivity, and a poorer adhesion in comparison with its counterpart on the glass substrate. Another disadvantage of the plastic substrate is its relatively high permeability of the moisture and oxygen that can affect the OLED electrical stability.5

In this letter we describe the optoelectrical characteristics of a high-efficiency, bilayer organic polymer light-emitting device with aluminum cathode fabricated on a thin, flexible plastic substrate. The device consists of a hole transporting (amine-fluorene) and an emissive (benzothiadiazole-fluorene) conjugated organic polymer layers. The best heterostructure device has a green light emission and a maximum luminance higher than 2000 cd/m². Device shows a maximum emission of ~56.2 cd/A and, accordingly, a maximum luminous and external quantum efficiency of ~9.0 lm/W and ~15%, respectively. This organic light-emitting diode performance is acceptable for flat panel display applications. © 2000 American Institute of Physics. [S0003-6951(00)01206-7]
where \( V'_0 = S'kT/q \) and \( V''_0 = S''kT/q \) are the fitting parameters; \( S' \) and \( S'' \) are the slopes of the \( I-V \) characteristics; \( k \) is Boltzmann constant, \( T \) (=300 K) is the temperature, \( q \) is the charge of electron, and \( I'_0 \) and \( I''_0 \) are constants. The \( V' \) and \( V'' \) represent the voltage drops across the two junctions connected in series, and \( R_e \) is the equivalent device series resistance. The best fit (solid line in Fig. 2) to the experimental data has been obtained for the following fitting parameters: \( I'_0 \approx 4.07 \times 10^{-9} \text{mA/cm}^2 \), \( I''_0 \approx 2.27 \times 10^{-2} \text{mA/cm}^2 \), \( V'_0 \approx 0.44 \), and \( V''_0 \approx 3.35 \).

The above equations describe the deep gap state (defects)-assisted multistep tunneling process\(^{10}\) that can take place across both junctions; and the \( V'_0 \) and \( V''_0 \) parameters represent the number of steps (involving a series of closely spaced deep gap defects) that carrier must travel through a given junction. These defects can be associated with the impurities and/or carbon or nitrogen point defects. In our case we assume that the HTL/EL junction is behaving like a \( p^+-p^- \) heterojunction, while the EL/Al junction is acting as a typical metal–semiconductor junction. Also the ITO/HTL (\( p^- \)-like) junction is assumed to be ohmic-like in our OLEDs, and therefore this contact will not provide any significant contribution to the \( I-V \) characteristic. Hence, at a constant electrical field the OLED carrier current is mainly limited by both junctions (e.g., multistep carrier tunneling/hopping through the HTL/EL junction and electron injection followed by the multistep tunneling/hopping through the EL/Al junction), rather than by the carrier transport through the polymer bulk; although the influence of the deep gap states-assisted bulk carrier transport can also be important in the OLED. In agreement with our experimental data the OLED injected current level (and the OLED brightness) is indeed higher if the low work function metals are used.\(^7\)

In the forward bias conditions, i.e., ITO is positively biased with respect to Al, the light emission occurred when the applied voltage is greater than \( \sim 6 \text{ V} \). No light emission is observed when the reversed bias is applied to the OLED structure, i.e., Al is positively biased with respect to ITO. In this case, the reverse resistance of the OLED is large. According to the Ohm’s law, a reverse voltage divided by a large device resistance yields a very small reverse current (e.g., \( 4.1 \times 10^{-4} \text{mA/cm}^2 \) at \( -25 \text{ V} \) in Fig. 2) that is insufficient to cause the light emission. A certain current density level (around \( 0.01 \text{ mA/cm}^2 \) in our case) is needed to produce the visible light. The electrons and holes that are injected, and transported through the junctions by multisteps process under the forward bias, into the polymer LUMO and HOMO bands are relaxed to negatively and positively charged polarons through electron-phonon lattice coupling. These polarons will move toward each other under the influence of the applied electrical field and they will recombine on a certain segment of polymer chain to form the singlet excitons. Through the resonance interactions, these singlet excitons can form the excimers, which will emit light Stokes shifted with respect to the absorption spectrum. The proposed mechanism of the light generation is in agreement with our experimental data.\(^{6,7}\)

Figure 3 shows a typical variation of the OLED brightness (or luminance, \( L \) in \( \text{cd/m}^2 \)) with the applied current density (the inset shows the OLED spectral distribution of luminance obtained at the different applied current densities). The integration of the OLED spectral distribution of luminance over the wavelength yields the luminance of the OLED. A near-linear relationship obtained between the luminance and the applied current density \( (J \text{ in mA/cm}^2) \), with a gradual saturation at \( J \approx 5 \text{ mA/cm}^2 \), can be described by the following equation:

\[
L \approx AJ + BJ^2, \tag{3}
\]

where \( A(354 \pm 28) \) and \( B(-14 \pm 3) \) are constants. At the low current injection this equation can be approximated by: \( L \propto J \); and at the high current injection the light-emission saturates at a certain level (in our case at \( \sim 2300 \text{ cd/m}^2 \) for
The inset shows the OLED photon density distribution spectra for different applied current densities.

$J \approx 12.6 \text{ mA/cm}^2$. The saturation in this curve is due to the OLED heating that is usually observed at the high current injection levels.\(^{12}\)

The evolution of the OLED external quantum efficiency (number of the photons emitted/electron supplied by the external circuit, e.g., applied current density) at different applied current densities is shown in Fig. 4 (the inset illustrates the photon density distribution spectra at the different applied current densities). The integration of the areas under the photon density distribution provides the number of photons emitted per unit area. We can conclude from this figure that the OLED emission efficiency increases rapidly at first then drops down gradually and starts slowly to flat out at about 4.78 mA/cm\(^2\) (27.7 cd/A) with the increasing applied current density (or luminance). The evolution of device external quantum efficiency also follows the same trend with the applied current density, and the maximum emission efficiency, luminous, and external quantum efficiencies for our OLED are $\sim 56.2 \text{ cd/A}$, $\sim 8.96 \text{ lm/W}$, and $\sim 14.8\%$, respectively. This excellent OLED performance could be obtained through the development of the following technologies for this device: minimum ITO surface roughness [root-mean-square (rms) $\sim 2.85 \text{ nm}$ measured over $\sim 30 \times 30 \mu\text{m}^2$]; thin flexible substrate (~0.2 mm); adequate refractive index matching between polymer (~1.72) and substrate; good control of the polymer morphology (aggregated species) during spin coating through the optimization of the solvent concentration of the polymer solution and the spin speed; and adequate electron transport within the EL polymer through the introduction of benzothiadiazole group (electron-withdrawing group), e.g., EL polymer is used for both the electron transport and light emission. The reduction of the external quantum efficiency at the higher current density is due to the light-induced defects creation (bond breaking) and/or device internal field modification leading to reduction of the density of the excimer states.

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