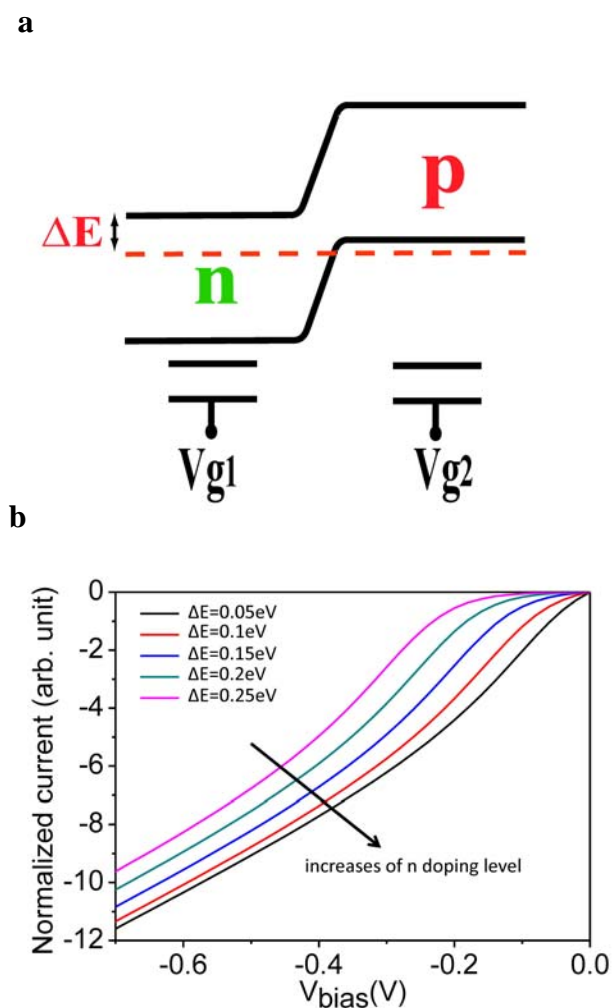


# A Fully Tunable Single-Walled Carbon Nanotube Diode

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## SUPPORTING INFORMATION

### A. Simulated $I$ - $V$ characteristics of the carbon nanotube tunneling junction under reverse bias voltage



**Figure S1 a**, Energy band diagram under equilibrium with fixed degenerated p doping and varying n doping, represented by changes in  $\Delta E$ . **b**, Simulated band-to-band tunneling in the reverse bias region.

Fig. S1(a) depicts band diagram under different equilibrium conditions. Fermi-level is initially near the band edge of the p side, and by changing  $V_{g1}$  gate voltage, doping level in the n side ( $\Delta E$ ) can be modulated. To numerically model  $I$ - $V$  characteristics under the reverse bias voltage, Fermi's golden rule is employed to calculate the band-to-band tunneling current<sup>1</sup>:

$$I_{tunneling} \propto \int_{E_{cn}}^{E_{vp}} (F_c(E) - F_v(E)) T N_c(E) N_v(E) dE$$

Here,  $F_c(E)$  and  $F_v(E)$  are Fermi-Dirac distribution functions at room temperature, and  $T$  is the tunneling probability through the junction potential barrier.  $N_c(E)$  and  $N_v(E)$  are density-of-states (DOS) of carbon nanotube for the conduction band and valence band, respectively. A large number of electronic states are near the band edge due to van Hove singularities. The equation integrates from the conduction band edge in the n side ( $E_{cn}$ ) to the valence band edge in the p side ( $E_{vp}$ ). From the equation, it is obvious that tunneling current is mainly decided by the overlap integral of DOS ( $\int N_c(E) N_v(E)$ ), modulated by Fermi energy difference ( $\int F_c(E) - F_v(E)$ ) between p and n side. For simplicity, only interband tunneling within the first subband is taken into consideration. Tunneling within higher subbands is negligible due to relative low tunneling probability<sup>2</sup>. The simulation results are plotted in the Fig. S1(b). Each curve shows similar  $I$ - $V$  characteristics, with a nearly exponential current increase due to a large of DOS at van Hove singularity, and then followed by linear dependence under reverse bias voltage. In addition,  $I$ - $V$  curve shifts toward downward direction by increasing doping level in the n side, which shows the same dependence as our experimental curves in the Fig. 3a. This theoretical simulation provides quantitative and predictive understanding of SWNT tunneling junction.

### References for Supplementary Information

- 1 S. M. Sze, and K. K. Ng, Physics of Semiconductor Devices, Wiley, New York, 2006.
- 2 Jena, D.; Fang, T.; Zhang, Q.; Xing, H. L. Appl. Phys. Lett. **2008**, 93, 112106.